

Feedback from last 2 lectures

- Bigger, better red boxes!
 Too much writing per slide
 More small boxes (equations etc) than large ones.
 - Maybe too much text more description, less reading out?
 - More space in handouts for notes / revising
 - Anything else?

Summary from last time:

- formal definitions of cross section
- definitions of Rutherford and Mott cross sections for coulomb scattering
- diffraction effects in scattering, determination of charge/mass distribution

For this lecture:

Implicit assumptions so far: nuclei are spheres

What determines the shape of a nucleus?



<u>Answer:</u> the attractive interactions between components form a droplet in order to minimize the number of "high energy" components at the surface (similar to the attractive forces between molecules in a droplet of water)

The binding energy of a nucleus:

- the energy available to hold nucleus together

Think of it this way:

- Take bunch of well-separated nucleons: binding energy is zero

- Bring them together: strong force glues them together. However, energy has to come from somewhere: binding energy must come from a **reduction** in nuclear mass

Formally, it is the difference between mass of component protons and neutrons and that of actual nucleus, related through $E = mc^2$:

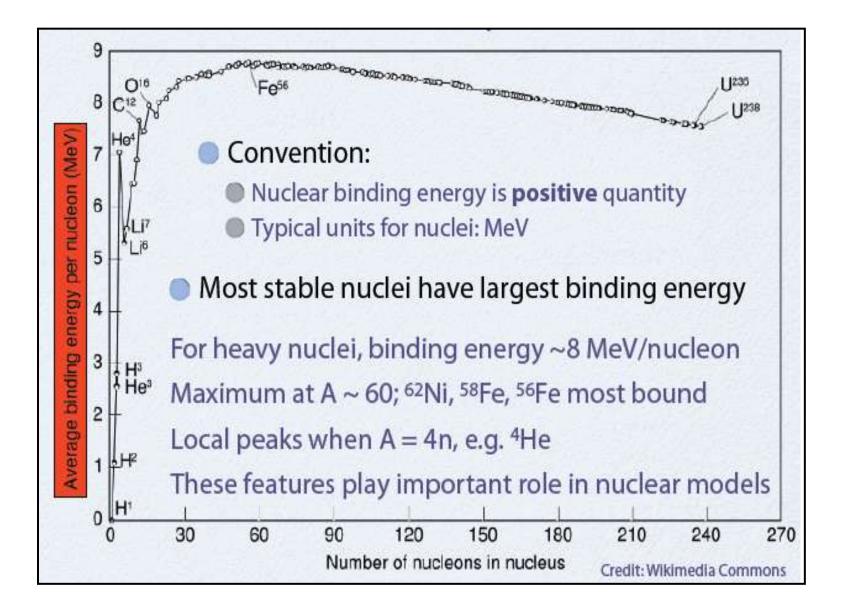
$$B(A,Z) = Z m_p c^2 + N m_n c^2 - M(A,Z) c^2$$

Binding energy is a **positive quantity**

(don't get confused here - the strong potential in which the nucleons sit is negative)

Binding energy per nucleon

- → the average energy state of nucleon is a sum of high energy "surface" nucleons with low energy "bulk" nucleons
- \rightarrow nucleus minimizes energy by minimizing surface area a sphere 4



The liquid drop model: a.k.a. "semi empirical mass formula" [SEMF]

Consider analogy with liquid drop:

- Liquids often considered as non-compressible
- Density constant, independent of radius

Then: radius **R** ~ **n**^{1/3} where n is number of molecules/nucleons in drop



Assume each molecule/nucleon bound with energy = -a (i.e. energy required to remove from drop against force due to all others)

B = an

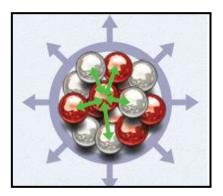
 \rightarrow Total binding energy given by:

However, only bound into drop on one side: Will see reduction in potential energy ~ $4\pi R^2 T$ i.e. proportional to surface area of drop and surface tension T

"Bulk" and "surface" terms should have <u>opposite</u> sign: absence of strong force on other side of surface makes nucleus <u>less</u> stable, <u>decreasing</u> binding energy

→ Using R~ $n^{1/3}$ and substituting $\beta = 4\pi$ T:

$$\mathsf{B} = \mathsf{an} - \beta \mathsf{n}^{2/3}$$

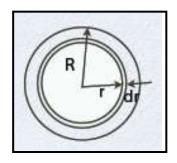


This describes basics of strong force interaction. However, proton – proton repulsion also important.

→ Including electrostatics in the liquid drop model:

In charged drop, repulsive force acts to unbind drop - reduces binding energy

Get change in binding energy by calculating the electrostatic potential of charge Q, distributed uniformly throughout.



How to do it: take uniform charge density ρ . Start with sphere radius r and charge $4/3\pi\rho r^3$. Then consider thin shell from r to r+dr with charge $4\pi\rho r^2$ dr. Calculate work done to bring shell from infinity to radius r, then integrate this from r=0 to R and express it in terms of total charge.

The result:

$$U = \frac{3Q^2}{20\pi\epsilon_0 R}$$

→ Again using R~n^{1/3} and substituting γ for all constants except charge: B = an - $\beta n^{2/3} - \gamma Q^2 n^{-1/3}$

All nucleons (A) carry strong force; only protons (Z) charged \rightarrow Substitute n \rightarrow A and Q \rightarrow Z in liquid drop model:

$$B(A, Z) = a_v A - a_s A^{2/3} - a_c Z^2 A^{-1/3}$$
Constants:
$$A_v \text{ for volume term}$$

$$a_s \text{ for surface term}$$

$$a_c \text{ for Coulomb term}_c$$

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Assumptions so far:

- 1) Nucleus is spherical
- 2) Nucleons behave like molecules in water drop:
 - → Short-range attractive force holding them together with shorter-range repulsive force stopping collapsing onto each other
 → Nuclear density is constant

Just including bulk, surface and electrostatic terms:

$$B(A, Z) = a_v A - a_s A^{2/3} - a_c Z^2 A^{-1/3}$$

However, formula inadequate as is:

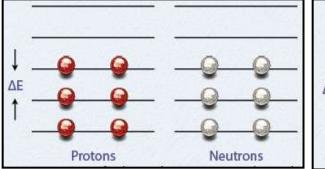
→ For fixed A, binding energy maximised when Z = 0Conversion of neutrons to protons allowed via β-decay, but $Z \rightarrow 0$ not seen in nature

→ Why?? Model not complete: two **non-classical terms** required:

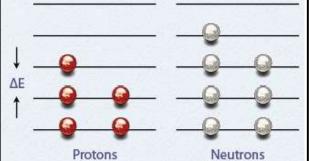
Both protons and neutrons are fermions - asymmetric wave functions and spin ¹/₂ Thus, they obey Pauli exclusion principle:

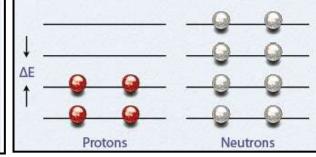
- Protons and neutrons must be arranged in separate energy levels: thus, consider two independent potential wells with identical energy levels, one for protons and one for neutrons
- Only two protons/neutrons possible in any given energy level, one with spin-up and the other with spin-down

Breaking symmetry between Z and N takes energy:



Most stable configuration





N–Z = 2: Reduced p-p repulsion But: moving proton has cost ΔE

N–Z = 4: has cost $2\Delta E$

For more states:

N-Z	Step energy (ΔE)	Cumulative energy (ΔE)
2	1	1
4	1	2
6	3	5
8	3	8
10	5	13
12	5	18
14	7	25

Cumulative energy change from lowest energy N = Z given by ~ $(N - Z)^2 / 8 \times \Delta E$

This can be rewritten as ~ $(A - 2Z)^2 / 8 \times \Delta E$

However, ΔE is not constant. From QM: energy levels for particle in 3D finite well follow $\Delta E \sim 1 / R^3$. Again using R ~ n^{1/3} = A^{1/3} $\rightarrow \Delta E \sim 1 / A$

Collecting constants together, change in

binding energy given by:

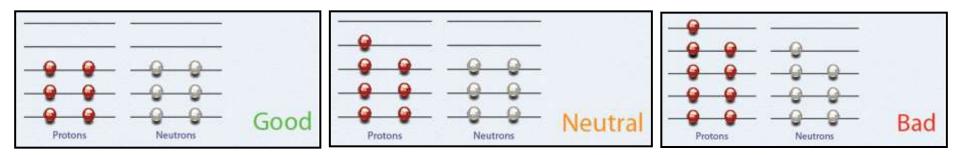
$$-a_a(A-2Z)^2 A^{-1}$$

i.e. asymmetry between protons and neutrons again **reduces** the binding energy, hence -ve

Spin pairing in the liquid drop model:

Spin pairing favours pairs of fermionic nucleons (similar to electrons in atoms) i.e. a pair with opposite spin have lower energy than pair with same spin

Best case: even numbers of both protons and neutrons Worst case: odd numbers of both protons and neutrons Intermediate cases: odd number of protons, even number of neutrons or vice versa



 \rightarrow Subtract small energy δ required to decouple nucleons from binding energy:

 $\begin{array}{rcl} \delta &=& +a_p A^{-1/2} & \mbox{ for both N \& Z odd} \\ &=& 0 & \mbox{ for N even, Z odd / Z even, N odd} \\ &=& -a_p A^{-1/2} & \mbox{ for both N \& Z even} \end{array}$

→ a_p collects constants, A^{-1/2} dependence provides best empirical fit to data

 \rightarrow subtracting δ <u>reduces</u> BE for N and Z both odd

 $\boldsymbol{\rightarrow}$ subtracting $\delta \;\underline{\textbf{adds}}$ small amount to BE for N and Z both even

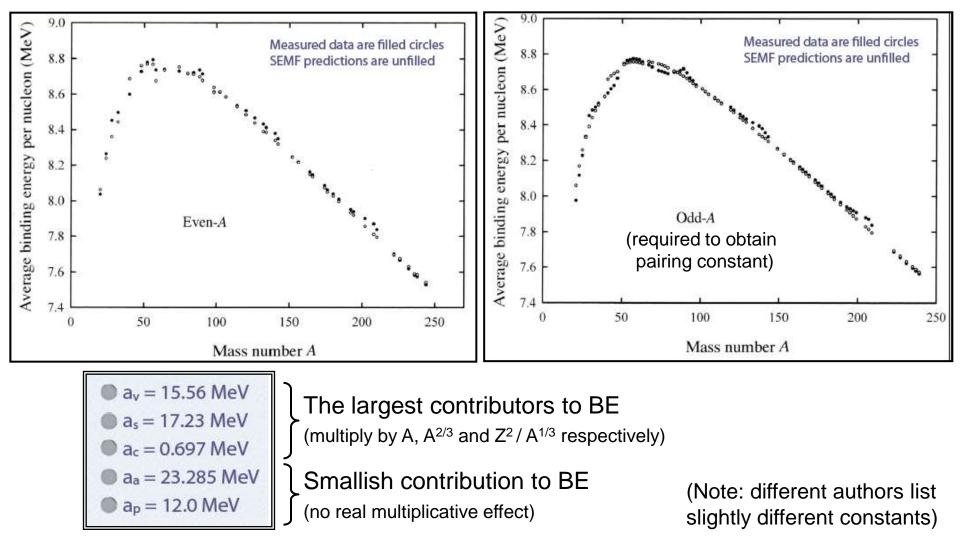
The constants:

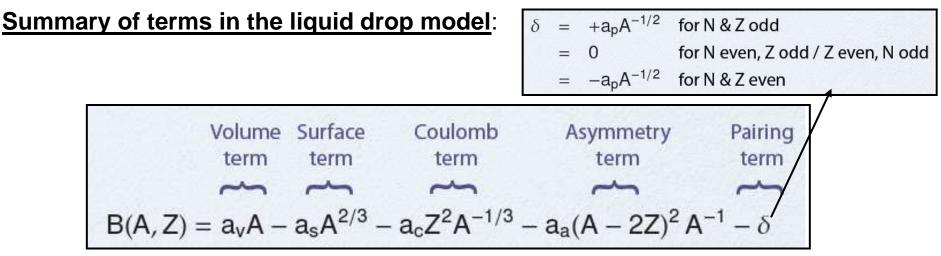
Why semi-empirical?

Constants not defined on basis of theory

Derived by fitting to measured masses

For A > 20, accuracy generally better than 0.1 MeV (i.e. error < 1% of total mass)





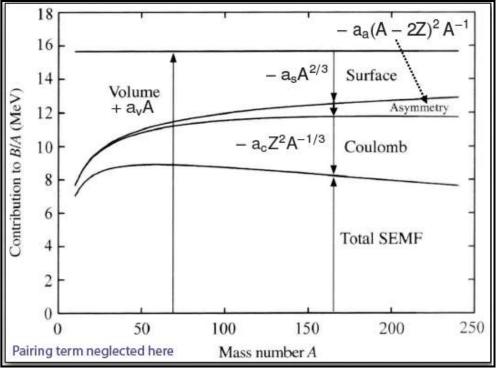
binding energy "comes" from volume term (i.e. the strong interaction between nucleons)
all other terms reduce binding energy:

Can also write in terms of mass:

 $M(A,Z) c^2 = Z m_p c^2 + N m_n c^2 - B(A,Z)$

(semi empirical mass formula)

Good prediction of nuclear masses: assumptions must be reasonable

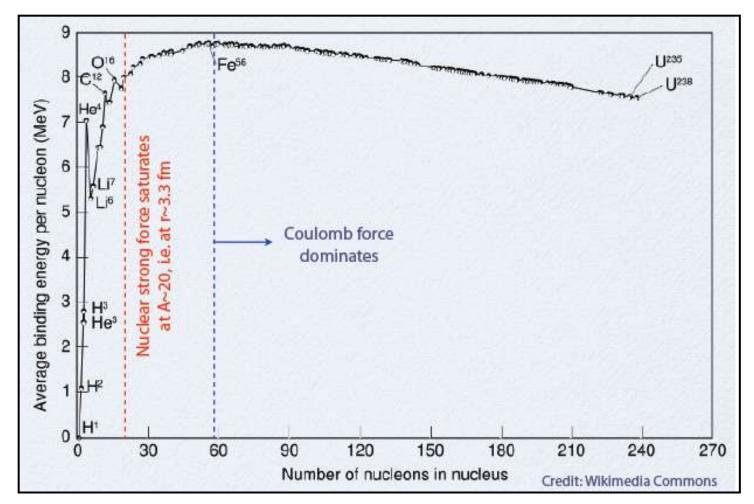


Implications of the liquid drop model for the nuclear force:

The attractive (volume) term is proportional to A, so strong force must be short range. If every nucleon interacted with every other nucleon, would go as $\sim A^2$ for large A.

- \rightarrow inter-nucleon force saturates very quickly on increasing A
- → must also become repulsive at small separations (prevents collapse of nucleus)

By contrast, electrostatic component goes as $\sim Z^2$ for large Z, i.e. does <u>not</u> saturate:



Typical exam questions:

1) Calculate the binding energy of uranium $^{238}_{\ 92} U$

A = 238 Z = 92 N = 238 - 92 = 146	Binding energy terms (in MeV): Volume: $a_vA = 15.56 \times 238 = 3703.28$ Surface: $a_sA^{2/3} = 17.23 \times (238)^{2/3} = 661.71$ Coulomb: $a_cZ^2A^{-1/3} = 0.697 \times (92)^2 \times (238)^{-1/3} = 951.95$ Asymmetry: $a_a(A-2Z)^2A^{-1} = 23.285 \times (238 - 184)^2 \times 238^{-1} = 285.29$ Pairing: $a_pA^{-1/2} = 12.0 \times (238)^{-1/2} = 0.78$
B = +Volume - Surface	- Coulomb - Asymmetry + Pairing (N, Z even)
= 3703.28 - 661.71 - 95	51.95 - 285.29 + 0.78

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= 1805.10 MeV
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 \rightarrow correct to 0.19%

2) What percentage is the binding energy of the total nuclear mass? Masses of unbound nucleons: p = 938.27 MeV/c², n = 939.57 MeV/c² Protons: 92 x 938.27 = 86 320.84 MeV/c2

Neutrons: 146 x 939.57 = 137 177.22 MeV/c2

Total: = 223 498.06 MeV/c2

Mass of bound nucleus:

= Mass of unbound nucleons - Binding energy = 223 498.06 - 1805.10 = 221 692.96 MeV/c²

 \rightarrow BE is < 1% of total mass energy of nucleus

Another example – Energy given out by nuclear reaction:

Calculate the energy released when an alpha particle is emitted by ²³⁵U nucleus.

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The reaction: {}^{235}U \rightarrow \alpha + {}^{231}Th
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Need to know binding energies of all components in reaction:

²³⁵U : B(235,92) = 3642.5 - 639.766 - 976.8 - 254.566 = 1771.368 MeV

²³¹Th : B(231,90) = 3580.5 - 632.486 - 939.928 - 258.974 = 1749.113 MeV

However, mass of alpha particle not well predicted by SEMF

- it is a particularly stable "magic" nucleus (see later)
- in practice, simply need to be given it!
- \rightarrow BE_{α} = 28.3 MeV/c²

Energy released = BE_{α} + B(A-4,Z-2) - B(A,Z)

= 28.3 + 1749.1 - 1771.4 = 6.0 MeV

 \rightarrow This energy is released as kinetic energy of α particle

Summary:

- explain and understand each of the terms in the liquid drop model
- be able to calculate nuclear masses and binding energies

For next time:

Limitations of the liquid drop model